

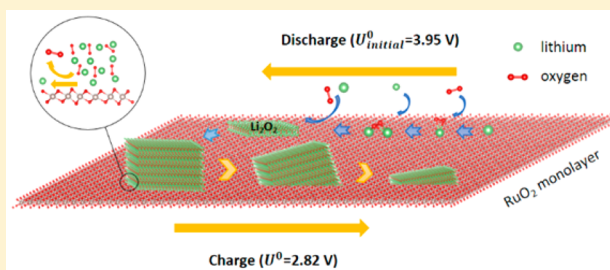
# RuO<sub>2</sub> Monolayer: A Promising Bifunctional Catalytic Material for Nonaqueous Lithium–Oxygen Batteries

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## S Supporting Information

**ABSTRACT:** Rutile RuO<sub>2</sub> has been widely regarded as an excellent catalyst for the oxygen reduction reaction (ORR) and oxygen evolution reaction (OER) in nonaqueous lithium–oxygen batteries and achieved superior performance, but the catalytic activity of RuO<sub>2</sub>'s polymorph, RuO<sub>2</sub> monolayer, has been less studied. In this work, we study the catalytic activities of both rutile RuO<sub>2</sub> and RuO<sub>2</sub> monolayer for ORR and OER in the battery using density functional theory method. Computational results show that the RuO<sub>2</sub> monolayer exhibits a higher catalytic activity than the rutile RuO<sub>2</sub> does. More interestingly, it is found that during discharge a similar lattice structure between RuO<sub>2</sub> monolayer and Li<sub>2</sub>O<sub>2</sub> {0001} surface can induce the formation of crystallized Li<sub>2</sub>O<sub>2</sub> with the conductive {0001} surface exposed, whereas during charge the RuO<sub>2</sub> monolayer can attract the remaining Li<sub>2</sub>O<sub>2</sub> to its surface spontaneously, thus maintaining the solid–solid reaction interface. Our results suggest that the RuO<sub>2</sub> monolayer is a promising catalytic material for nonaqueous lithium–oxygen batteries.



## 1. INTRODUCTION

Nonaqueous lithium–oxygen batteries have attracted increasing attention because of their superhigh specific capacities;<sup>1–7</sup> however, as an infant stage technology, this novel energy-storage system is still suffering from several severe issues, such as the high charge overpotential, sluggish reaction kinetics, and poor cycling stability, which seriously hindered its commercialization. A common strategy to solve these issues is using catalyst to lower the charge overpotential and accelerate the reactions, which could also alleviate the decomposition of electrolyte and cathode materials, thus enhancing cycling stability. Many possible catalysts have been tested,<sup>8–14</sup> among which rutile RuO<sub>2</sub> is a popular candidate for its superior catalytic activity and good chemical stability in the highly oxidative environment of nonaqueous lithium–oxygen batteries.<sup>15–17</sup> As the polymorph of rutile RuO<sub>2</sub>, RuO<sub>2</sub> monolayer exfoliated from lamellar ruthenates is also conductive and has been used as super capacitor and catalyst for the oxygen reduction reaction (ORR) in aqueous system.<sup>18–21</sup> Being a 2D metal oxide, RuO<sub>2</sub> monolayer can achieve a large surface-to-mass ratio and exist stably during the operation of nonaqueous lithium–oxygen batteries. It is also expected to catalyze the reactions uniformly without the necessity to introduce defects or functional groups as carbon material does, making it a possible candidate for the catalyst or catalytic cathode of nonaqueous lithium–oxygen batteries. Recently, Liao et al.<sup>22</sup> used lamellar ruthenate as a precursor to obtain RuO<sub>2</sub> nanosheet to be a catalytic cathode of nonaqueous lithium–oxygen batteries and achieved excellent performance. In their experiment, however, a considerable amount of exfoliated RuO<sub>2</sub>

monolayer went through phase change and turned to rutile RuO<sub>2</sub> during the heat treatment process. Thus, whether RuO<sub>2</sub> monolayer contributed to the excellent performance is still unclear. Further detailed computational and experimental investigations are needed to identify whether RuO<sub>2</sub> monolayer is a promising effective catalytic material for nonaqueous lithium–oxygen batteries.

In this work, using density functional theory (DFT) method, we investigated the catalytic activities of both rutile RuO<sub>2</sub> and RuO<sub>2</sub> monolayer for the ORR and oxygen evolution reaction (OER) in nonaqueous lithium–oxygen batteries. The dominant surface of rutile RuO<sub>2</sub> was identified according to the Wulff construction to be representative in the study of the catalytic activity of rutile RuO<sub>2</sub>. The adsorption behaviors of the lithium ion, intermediate discharge product LiO<sub>2</sub>, and discharge product Li<sub>2</sub>O<sub>2</sub> on both RuO<sub>2</sub> monolayer and the dominant surface of rutile RuO<sub>2</sub> were investigated to study the initial discharge process. The electronic properties of RuO<sub>2</sub> monolayer before and after the deposition of different layers of Li<sub>2</sub>O<sub>2</sub> were obtained to study the following discharge process when Li<sub>2</sub>O<sub>2</sub> accumulates after the initial deposition. Interfacial models among RuO<sub>2</sub> monolayer or the dominant surface of rutile RuO<sub>2</sub>, Li<sub>2</sub>O<sub>2</sub>, and electrolyte were built to study the charge process.

## 2. COMPUTATIONAL DETAILS

DFT calculations were performed using the ABINIT<sup>23–25</sup> software package with the projector-augmented-wave (PAW)

Received: January 4, 2016

Revised: March 3, 2016

Published: March 9, 2016

method<sup>26</sup> and Perdew–Burke–Ernzerhof (PBE) generalized gradient approximation (GGA).<sup>27</sup> The cutoff energy for the plane-wave basis was 20 Ha. Monkhorst–Pack scheme<sup>28</sup> was used for the k-point sampling, and the spacing of the k-point mesh was set to be  $<0.05 \text{ \AA}^{-1}$ . All atoms were relaxed to a force tolerance of  $0.02 \text{ eV \AA}^{-1}$  or less for the geometric optimization. The Bader charge partition analysis<sup>29,30</sup> implemented in the ABINIT software was employed to study the related charge transfer. The density of states (DOS) were calculated using Heyd–Scuseria–Ernzerhof (HSE) hybrid density functional<sup>31,32</sup> implemented in Quantum Espresso package.<sup>33</sup> One-quarter ( $\alpha = 0.25$ ) of the local DFT exchange was replaced by the unscreened and nonlocal Fock exchange.

The crystal structure of  $\text{RuO}_2$  monolayer was obtained by modifying the crystal structure of  $\text{Na}_x\text{RuO}_2 \cdot y\text{H}_2\text{O}$ .<sup>34</sup> The optimized lattice parameters of  $\text{RuO}_2$  monolayer (listed in Table S1) agree well with experimental reports.<sup>19</sup> The crystal structure of rutile  $\text{RuO}_2$  and  $\text{Li}_2\text{O}_2$  was taken from experimental results.<sup>35,36</sup>

Slab model adding vacuum layers was used to calculate the surface energies. All slabs are symmetrized and contain more than seven repeating layers with a vacuum layer thicker than  $10 \text{ \AA}$  to achieve convergence within  $1 \text{ meV \AA}^{-2}$  for the surface energies. The surface energies were calculated by

$$\gamma = \frac{1}{2A} [G_{\text{slab}} - N_{\text{Ru}}\mu_{\text{Ru}}^{\text{bulk}} - N_{\text{O}}\mu_{\text{O}}^{\text{bulk}}] \quad (1)$$

where  $A$  is the slab surface area,  $G_{\text{slab}}$  is the total free energy of the slab supercell,  $N_{\text{Ru}}$  and  $N_{\text{O}}$  are the numbers of ruthenium and oxygen atoms, and  $\mu_{\text{Ru}}^{\text{bulk}}$  and  $\mu_{\text{O}}^{\text{bulk}}$  are the chemical potentials of ruthenium and oxygen atoms, respectively. The chemical potentials of  $\text{RuO}_2$  are correlated by

$$\mu_{\text{Ru}}^{\text{bulk}} + 2\mu_{\text{O}}^{\text{bulk}} = G_{\text{RuO}_2} \quad (2)$$

Thus the surface energies can be rewritten as

$$\gamma = \frac{1}{2A} [G_{\text{slab}} - N_{\text{Ru}}G_{\text{RuO}_2} + (2N_{\text{Ru}} - N_{\text{O}})\mu_{\text{O}}^{\text{bulk}}] \quad (3)$$

The Wulff shape was constructed using the WinXmorph software.<sup>37</sup>

A  $3 \times 3$  supercell of  $\text{RuO}_2$  monolayer and a  $2 \times 2$  supercell of rutile  $\text{RuO}_2$  {001} surface with a  $15 \text{ \AA}$  vacuum layer were used to study the adsorption behavior. For the rutile  $\text{RuO}_2$  {001} surface, the two exposed surfaces are kept symmetrically equivalent for all calculations to eliminate the polarization effect. For  $\text{LiO}_2$  and  $\text{Li}_2\text{O}_2$ , the adsorption energies were calculated using

$$E_{\text{ads}} = E_{\text{tot}}(S + A) - E_{\text{tot}}(A) - E_{\text{tot}}(S) \quad (4)$$

where  $E_{\text{tot}}(S+A)$  is the total energy of the adsorbed substrate,  $E_{\text{tot}}(A)$  is the total energy of adsorbate, and  $E_{\text{tot}}(S)$  is the total energy of substrate. For  $\text{Li}^+ + \text{e}^-$ , the adsorption energy was calculated by

$$E_{\text{ads}} = E_{\text{tot}}(S + A) - E_{\text{tot}}(S) - \mu_{\text{Li}} \quad (5)$$

where  $\mu_{\text{Li}}$  is equal to the chemical potential of bulk lithium.

To address the well-known oxygen overbinding issue in DFT calculations, we calculated the enthalpy of oxygen molecule using<sup>38</sup>

$$H(T = 0 \text{ K}, \text{O}_2) = 2H(T = 0 \text{ K}, \text{O}) - \Delta E^{\text{exptl}} \quad (6)$$

where  $H(T = 0 \text{ K})$  is the calculated zero-point energy of oxygen atom or molecule and  $\Delta E^{\text{exptl}}$  ( $5.12 \text{ eV}$ )<sup>38</sup> is the binding energy of oxygen from experiments. The chemical potential of oxygen was calculated by

$$\mu_{\text{O}_2}(T, P_{\text{O}_2}) = H_{\text{O}_2}(0 \text{ K}) + \Delta H_{\text{O}_2}(T) - TS_{\text{O}_2}^{\text{exptl}}(T) + k_{\text{B}}T \ln \left( \frac{P_{\text{O}_2}}{P_{\text{O}_2}^0} \right) \quad (7)$$

where  $\Delta H_{\text{O}_2}(T)$  is the enthalpy energy change from  $0 \text{ K}$  to  $T$ , for which we used diatomic ideal gas approximation as  $7/2k_{\text{B}}T$ , and  $S_{\text{O}_2}^{\text{exptl}}(T)$  is the entropy of oxygen at  $1 \text{ atm}$  and different temperatures obtained from experiments.<sup>39</sup>  $P_{\text{O}_2}^0$  is set to be  $1 \text{ atm}$ .

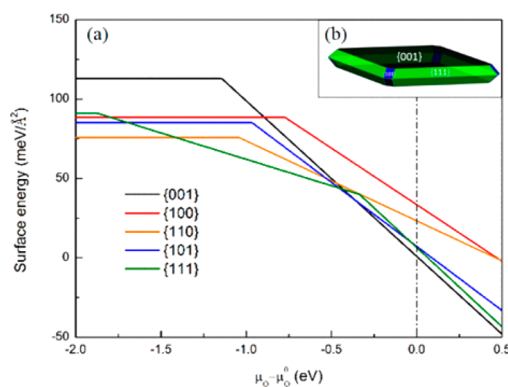
The energy profiles of the ORR/OER process were calculated by adding/removing a lithium atom or an oxygen molecule at each step. The reaction free energy of intermediate steps was calculated by

$$\Delta G = E - E_0 + \Delta N_{\text{Li}}(\mu_{\text{Li}} - eU) + \Delta N_{\text{O}_2}\mu_{\text{O}_2} \quad (8)$$

where  $E$  is the total energy of the considered slab model,  $E_0$  is the total energy of the initial slab model,  $\Delta N_{\text{Li}}$  and  $\Delta N_{\text{O}_2}$  are the numbers of lithium atoms and oxygen molecules added/removed for each step, and  $\mu_{\text{Li}}$  and  $\mu_{\text{O}_2}$  are the chemical potentials of lithium bulk and oxygen, respectively. The  $eU$  term was added to account for the electronic energy under applied potential  $U$ . The overpotential was defined by shifting all intermediates to  $\Delta G < 0$ , which is consistent with previous works.<sup>40–42</sup>

### 3. RESULTS AND DISCUSSION

**3.1. Representative Surface of Rutile  $\text{RuO}_2$ .** To identify the stable surfaces of rutile  $\text{RuO}_2$  in the operation condition of nonaqueous lithium–oxygen batteries, we calculated the surface energies of {001}, {100}, {101}, {110}, and {111} surfaces with different terminations (as shown in Figure S1). The calculated surface energies under different oxygen chemical potentials are shown in Figure 1a. Because nonaqueous lithium–oxygen

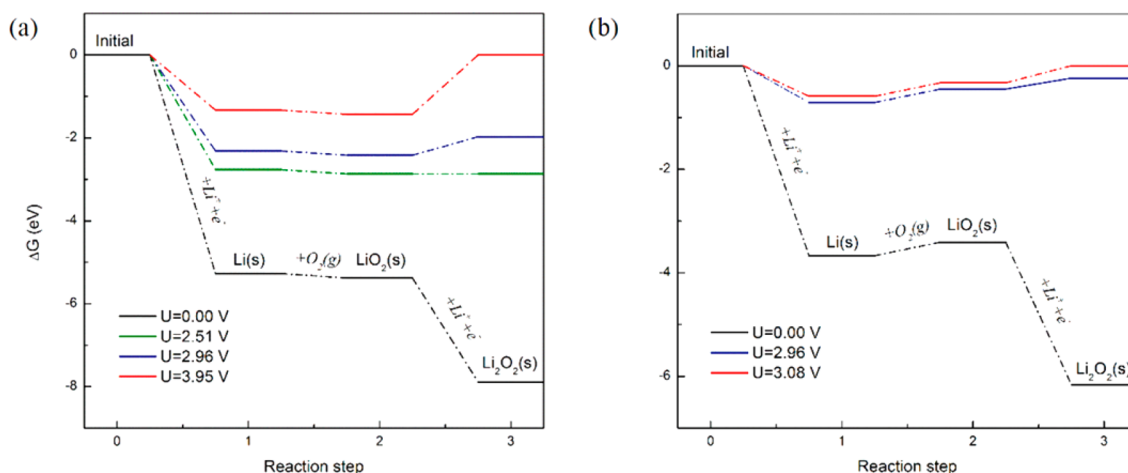


**Figure 1.** (a) Surface energies of rutile  $\text{RuO}_2$  at different oxygen chemical potentials and (b) the Wulff construction of rutile  $\text{RuO}_2$  when  $P(\text{O}_2) = 1 \text{ atm}$ .

batteries operate in an oxygen-rich condition, here we take the surface energies when  $P_{\text{O}_2} = 1 \text{ atm}$  to construct the Wulff structure of rutile  $\text{RuO}_2$ . From the Wulff shape shown in Figure 1b, it can be found that the {001} surface occupies most of the

Table 1. Adsorption of  $\text{Li}^+ + \text{e}^-$ ,  $\text{LiO}_2$ , and  $\text{Li}_2\text{O}_2$  onto the  $\text{RuO}_2$  Monolayer and Rutile  $\text{RuO}_2$  {001} Surface, Respectively

adsorbate	$E_{\text{ads}}$ (eV)		$D_{\text{O-O}}$ (Å)		$\delta(\text{e})$ of adsorbate	
	$\text{RuO}_2$ monolayer	$\text{RuO}_2\{001\}$	$\text{RuO}_2$ monolayer	$\text{RuO}_2\{001\}$	$\text{RuO}_2$ monolayer	$\text{RuO}_2\{001\}$
$\text{Li}^+ + \text{e}^-$	−5.27	−3.67	—	—	+0.91	+0.88
$\text{LiO}_2$	−4.18	−2.21	1.27	1.24	+0.55	+0.78
$\text{Li}_2\text{O}_2$	−5.02	−3.29	1.37	1.27	+0.91	+1.44

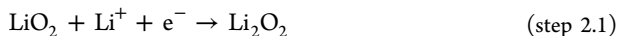
Figure 2. Energy profiles for the initial discharge process happened on (a) the surface of  $\text{RuO}_2$  monolayer and (b) rutile  $\text{RuO}_2$  {001} surface.

exposed surface area. Thus, in the following, we will take the {001} surface as a representative of rutile  $\text{RuO}_2$  to study its catalytic activities toward the ORR and OER in nonaqueous lithium–oxygen batteries.

**3.2. Initial Discharge Process on  $\text{RuO}_2$  {001} and  $\text{RuO}_2$  Monolayer.** It is now widely agreed that the ORR in nonaqueous lithium–oxygen batteries goes through a two-step reaction.<sup>1–6</sup> In the first step, the oxygen molecule will complex with a lithium ion and get an electron to form lithium superoxide as



Then, the formed  $\text{LiO}_2$  could either go through an electrochemical reaction route and be further reduced to  $\text{Li}_2\text{O}_2$  as



or go through a chemical disproportionation reaction route as



The electrochemical reactions take place only on the catalytic surfaces, whereas the chemical disproportionation reaction can happen anywhere. It is commonly believed that the large compact toroid  $\text{Li}_2\text{O}_2$  particles resulted from the chemical disproportionation reactions that happened in electrolyte.<sup>43,44</sup> From our calculation results on the adsorption behaviors of  $\text{Li}^+ + \text{e}^-$ ,  $\text{LiO}_2$ , and  $\text{Li}_2\text{O}_2$  onto both  $\text{RuO}_2$  monolayer and rutile  $\text{RuO}_2$  {001} surface, as listed in Table 1, it can be found that in the both cases the adsorption energies of  $\text{LiO}_2$  are quite large. The large adsorption energies of  $\text{LiO}_2$  could confine all reactions onto the surfaces in the initial stage of discharge process, rather than letting the  $\text{LiO}_2$  dissolve into the electrolyte and disproportionate into  $\text{Li}_2\text{O}_2$ , thus effectively suppressing the formation of large compact toroid particles that are hard to be charged back.<sup>45</sup> So, in this study, we only consider the disproportionation reactions that happened on the surfaces.

The detailed reaction route for the disproportionation reaction occurring on the surfaces can be much more complex. Here we consider only the reaction enthalpy change for an initial evaluation by defining the formation heat following Geng et al.'s<sup>45</sup> approach as

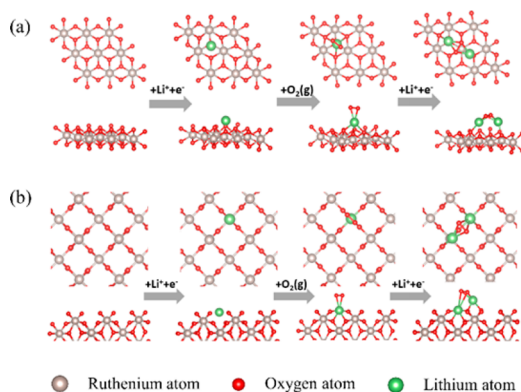
$$E_{\text{dis}} = E_{\text{Li}_2\text{O}_2@\text{surf}} - 2E_{\text{LiO}_2@\text{surf}} + E_{\text{surf}} + E_{\text{O}_2} \quad (9)$$

where  $E_{\text{O}_2}$ ,  $E_{\text{surf}}$ ,  $E_{\text{LiO}_2@\text{surf}}$  and  $E_{\text{Li}_2\text{O}_2@\text{surf}}$  represent the energy of oxygen molecule, considered surface, considered surface adsorbed with  $\text{LiO}_2$ , and considered surface adsorbed with  $\text{Li}_2\text{O}_2$ , respectively. The calculated  $E_{\text{dis}}$  for the disproportionation reactions that happened on the  $\text{RuO}_2$  monolayer is +3.17 eV, and for the rutile  $\text{RuO}_2$  {001} surface it is +0.953 eV, which means the disproportionation reactions on both surfaces are endothermic and can hardly be realized.

We now discuss the electrochemical route. As for both the  $\text{RuO}_2$  monolayer and rutile  $\text{RuO}_2$  {001}, all ruthenium atoms are coordinatively saturated, making the initial adsorption of oxygen not preferred, which can be further confirmed by the bond lengths of adsorbed oxygen molecules onto the surfaces, as shown in Figures S3 and S4. Thus, we study the electrochemical ORR in nonaqueous lithium–oxygen batteries by taking the reaction route  $\text{Li} \rightarrow \text{LiO}_2 \rightarrow \text{Li}_2\text{O}_2$ . The corresponding energy profiles for this process on different surfaces are shown in Figure 2. For both  $\text{RuO}_2$  monolayer and rutile  $\text{RuO}_2$  {001} surface, the ORR following the electrochemical reaction route are thermodynamically favored. Figure 2a shows the energy profiles of the ORR on the  $\text{RuO}_2$  monolayer at an open-circuit potential ( $U = 0$  V), the highest voltage at which discharging is energetically downhill for all steps ( $U = 2.51$  V), the equilibrium voltage of bulk  $\text{Li}_2\text{O}_2$  ( $U = 2.96$  V),<sup>39</sup> and the equilibrium voltage for the initial discharge process, which could just keep all of the intermediates having a negative energy ( $U = 3.95$  V). Figure 2b presents the energy profiles of the ORR on the rutile  $\text{RuO}_2$  {001} surface at an open-circuit potential ( $U = 0$  V), the equilibrium voltage of



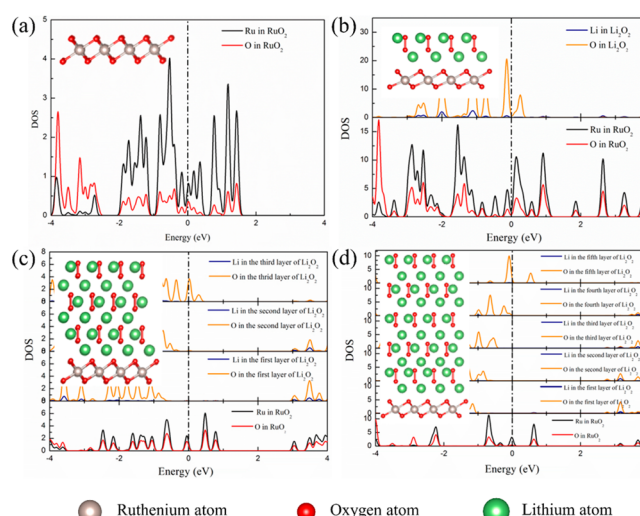
bulk  $\text{Li}_2\text{O}_2$  ( $U = 2.96$  V),<sup>39</sup> and the equilibrium voltage for the initial discharge process, which could just keep all of the intermediates having a negative energy ( $U = 3.08$  V). It can be found that the equilibrium voltage for the electrochemical ORR on  $\text{RuO}_2$  monolayer is as high as 3.95 V, much higher than that of the equilibrium voltage on rutile  $\text{RuO}_2$  {001} (3.08 V), indicating a higher catalytic activity of  $\text{RuO}_2$  monolayer toward the electrochemical ORR process. This result may also explain the experimental discharge curve of Liao et al., where an initial discharge voltage higher than 3.5 V was observed (corresponding to the equilibrium voltage as high as 3.95 V in the initial stage of the discharge process) and gradually drop down to a voltage plateau below 2.96 V (corresponding to the equilibrium voltage of the following discharge process that happened on  $\text{Li}_2\text{O}_2$  surface after  $\text{RuO}_2$  monolayer surface was fully covered by discharge product).<sup>22</sup> The reverse process of the initial discharge process is the charge process in its final stage, when there is only a very small number of  $\text{Li}_2\text{O}_2$  molecules adsorbed onto the surfaces waiting to be removed, corresponding to a high ending charge voltage for  $\text{RuO}_2$  monolayer, which also agrees with the experimental observation.<sup>22</sup> The geometry change along with the initial discharge process on the surfaces is shown in Figure 3.



**Figure 3.** Geometries for the initial discharge process happened on (a) the surface of  $\text{RuO}_2$  monolayer and (b) rutile  $\text{RuO}_2$  {001} surface.

**3.3. Electronic Conductance of  $\text{RuO}_2$  Monolayer Deposited with  $\text{Li}_2\text{O}_2$ .** As the discharge goes on, the active surfaces will be covered by  $\text{Li}_2\text{O}_2$  and show little influence toward the following discharge process (shown in Figures S5 and S6). From the results of our previous work,<sup>46</sup> the surface of  $\text{Li}_2\text{O}_2$  also has a strong adsorption ability toward  $\text{LiO}_2$ , and thus the following ORR reactions will mainly take place at the surfaces of  $\text{Li}_2\text{O}_2$ . Because  $\text{Li}_2\text{O}_2$  bulk is insulator, the electronic conductivity of  $\text{Li}_2\text{O}_2$  will become the main limiting factor for the following ORR.<sup>1–6</sup> For rutile  $\text{RuO}_2$ , its lattice structure and lattice parameters are quite different from those of  $\text{Li}_2\text{O}_2$ , which could induce the formation of amorphous  $\text{Li}_2\text{O}_2$  and increase the electronic conductivity through the formation of grain boundaries.<sup>15,45</sup>

For  $\text{RuO}_2$  monolayer, it should be noted that the crystal structures of  $\text{Li}_2\text{O}_2$  {0001} surface and  $\text{RuO}_2$  monolayer are similar, and their lattice parameters are quite close, as listed in Table S1. Considering that lithium atoms tend to adsorb onto the hcp hollows formed by the Ru atoms (shown in Figure 3a) and {0001} surface is the dominant surface of  $\text{Li}_2\text{O}_2$  in its equilibrium Wulff structure,<sup>38,42,47</sup> we propose that after the initial deposition process the formed  $\text{Li}_2\text{O}_2$  is very likely to adsorb onto the surface of  $\text{RuO}_2$  monolayer in the way shown



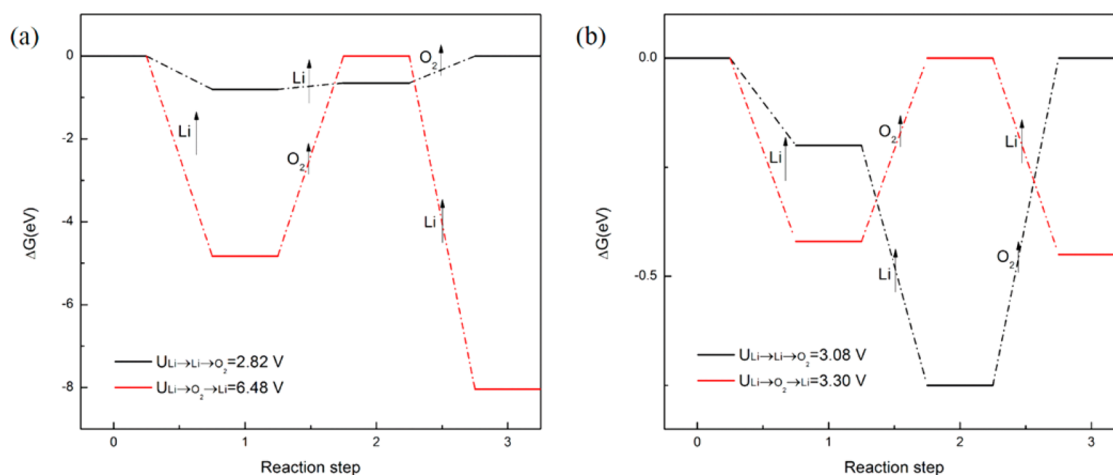
**Figure 4.** Density of states (DOS) map for the  $\text{RuO}_2$  monolayer before (a) and after deposition of one (b), three (c), and five (d) layers of  $\text{Li}_2\text{O}_2$ .

in Figure 4. In this way, the optimized lattice parameters of  $\text{Li}_2\text{O}_2$  after adsorbed onto  $\text{RuO}_2$  monolayer are quite close to that of bulk  $\text{Li}_2\text{O}_2$  (as shown in Figure S7).

We now present the study on electronic conductivity of  $\text{RuO}_2$  before and after the deposition of one, three, and five layers of  $\text{Li}_2\text{O}_2$ . HSE hybrid density functional was used to get accurate band information.<sup>31,32</sup> The constructed models and calculated DOS maps are shown in Figure 4. According to the calculation results,  $\text{RuO}_2$  monolayer is a metallic material (Figure 4a), which is consistent with experimental observations.<sup>18–21</sup> After the deposition of one layer of  $\text{Li}_2\text{O}_2$ , the  $\text{RuO}_2$  monolayer and the deposited  $\text{Li}_2\text{O}_2$  are both conductive (Figure 4b). If we continue to increase the  $\text{Li}_2\text{O}_2$  deposition (Figures 4c,d), the  $\text{RuO}_2$  monolayer will still be conductive, while the conductive  $\text{Li}_2\text{O}_2$  layer will shift up to the top layer exposed, which will be further away from  $\text{RuO}_2$ . At the same time, the  $\text{Li}_2\text{O}_2$  layer in contact with the  $\text{RuO}_2$  monolayer will exhibit to be an insulator with a band gap larger than 4.0 V, which is similar to that of bulk  $\text{Li}_2\text{O}_2$ .<sup>48,49</sup> This phenomenon is consistent with the previous result that the {0001} surface of  $\text{Li}_2\text{O}_2$  is conductive.<sup>58</sup> The surface conductivity would provide an electron pathway for the following discharge process.

Because a small lattice mismatch between  $\text{Li}_2\text{O}_2$  and  $\text{RuO}_2$  monolayer exists, the strain caused by lattice mismatch will become larger as the adsorbed  $\text{Li}_2\text{O}_2$  grows thicker.<sup>50</sup> This strain may induce the generation of defects and grain boundaries, which could further enhance the electronic conductivity.<sup>49,51</sup> The {111} surface of previously reported cathode material for nonaqueous lithium–oxygen batteries, TiC, also has a similar lattice structure and a small lattice mismatch compared with the  $\text{Li}_2\text{O}_2$  {0001} surface.<sup>52</sup> It is quite interesting to notice that the discharge product morphologies on TiC cathode<sup>53</sup> and  $\text{RuO}_2$  nanosheet cathode<sup>22</sup> are also similar; in both cases, a rare morphology of assembled thin discs was observed. We propose that this morphology may be caused by the stripping of  $\text{Li}_2\text{O}_2$  along the {0001} surface to relieve the internal strain, as shown in Figure S8, just as those observed in heteroepitaxial growth experiments.<sup>50,54,55</sup>

**3.4. Charge Process on Rutile  $\text{RuO}_2$  {001} and  $\text{RuO}_2$  Monolayer.** The involvement of solid-state discharge product makes the ORR and OER in nonaqueous lithium–oxygen

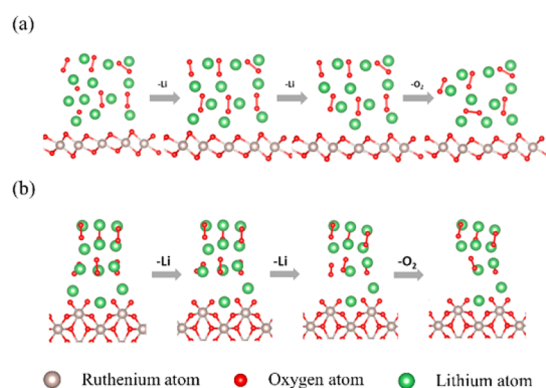


**Figure 5.** Energy profiles for the charge process of  $\text{Li}_2\text{O}_2$  happened on (a) the surface of  $\text{RuO}_2$  monolayer and (b) rutile  $\text{RuO}_2$  {001} surface following two possible reaction routes.

batteries asymmetric. As previously mentioned, the adsorption models built above to study the initial discharge process could also be used to interpret the end stage of the charge process but obviously not the beginning and middle stages. To describe the charge process more accurately, while keep the computational efforts reasonable, we built solid–solid interface models following Zhu et al.’s<sup>40,41</sup> approach. The same amount of nanoscale  $\text{Li}_2\text{O}_2$  was put onto both  $\text{RuO}_2$  monolayer and rutile  $\text{RuO}_2$  {001} surface, with {0001} surface of  $\text{Li}_2\text{O}_2$  in contact with them. After geometry optimization, for the rutile  $\text{RuO}_2$  {001} surface, the adsorbed  $\text{Li}_2\text{O}_2$  showed a tilt to fit the large lattice mismatch, as shown in Figure S9; for  $\text{RuO}_2$  monolayer, one of the oxygen–oxygen bonds near the interface was cleaved due to the lithium-rich termination.

According to previous studies,<sup>56</sup> the OER in nonaqueous lithium–oxygen batteries takes place at the three-phase interface of cathode/ $\text{Li}_2\text{O}_2/\text{O}_2$ . Thus, we chose a unit cell of  $\text{Li}_2\text{O}_2$  near the  $\text{RuO}_2$  monolayer or rutile  $\text{RuO}_2$  {001} surface/ $\text{Li}_2\text{O}_2$ /electrolyte interface to study the OER mechanism. Two possible reaction routes, namely,  $\text{Li}^+ \rightarrow \text{Li}^+ \rightarrow \text{O}_2$  (route 1) and  $\text{Li}^+ \rightarrow \text{O}_2 \rightarrow \text{Li}^+$  (route 2) were considered accordingly. The energy profiles for these two different routes are shown in Figure 5. For  $\text{RuO}_2$  monolayer, as shown in Figure 5a, the equilibrium voltage for route 1 is 2.82 V, much lower than that of 6.48 V for route 2. For rutile  $\text{RuO}_2$  {001} surface, as shown in Figure 5b, the equilibrium voltage for route 1 is 3.08 V, which is quite close to the charge voltages reported in experiments<sup>15,16</sup> and a bit lower than that of 3.30 V for route 2. Thus, route 1 will be the preferred reaction route in both situations. The equilibrium charge voltage for  $\text{RuO}_2$  monolayer is lower than that of rutile  $\text{RuO}_2$  {001} surface, indicating a higher catalytic activity during the initial and middle stages of the charge process.

The geometry changes of the adsorbed  $\text{Li}_2\text{O}_2$  along with the charge process following the route 1 are shown in Figure 6. It is quite interesting to find that for  $\text{RuO}_2$  monolayer, as shown in Figure 6a, after the removal of one unit cell of  $\text{Li}_2\text{O}_2$ , the remaining  $\text{Li}_2\text{O}_2$  will be attracted to move toward the  $\text{RuO}_2$  monolayer spontaneously (the center of mass of the adsorbed  $\text{Li}_2\text{O}_2$  moves toward the  $\text{RuO}_2$  layer about 0.2 Å), while similar behavior cannot be observed for the rutile  $\text{RuO}_2$  {001} surface, as shown in Figure 6b (the center of mass of the adsorbed  $\text{Li}_2\text{O}_2$  moves away from the  $\text{RuO}_2$  {001} surface about 0.4 Å).



**Figure 6.** Geometries for the charge process following route 1 happened on (a) the surface of  $\text{RuO}_2$  monolayer and (b) rutile  $\text{RuO}_2$  {001} surface.

We propose that this attraction effect could help maintain the solid–solid reaction interface until all adsorbed  $\text{Li}_2\text{O}_2$  is decomposed. Gittleston et al.<sup>57</sup> observed that for the gold cathode, after the beginning stage of the charge process when the first layer of  $\text{Li}_2\text{O}_2$  was consumed, the remaining  $\text{Li}_2\text{O}_2$  would be in poor contact with the cathode and can hardly be charged back. Our calculation results show that the vanishing of the reaction interface along with the charge process is unlikely to happen when the reactions take place on the surface of  $\text{RuO}_2$  monolayer.

#### 4. CONCLUSIONS

In summary, we comparatively investigated the ORR and OER in nonaqueous lithium–oxygen batteries occurring on both the surface of  $\text{RuO}_2$  monolayer and rutile  $\text{RuO}_2$ . The calculation results show that in nonaqueous lithium–oxygen batteries,  $\text{RuO}_2$  monolayer exhibits higher catalytic activity in both ORR and OER than rutile  $\text{RuO}_2$  does. In addition, during discharge, the similar lattice structure between  $\text{RuO}_2$  monolayer and  $\text{Li}_2\text{O}_2$  {0001} can induce the conductive  $\text{Li}_2\text{O}_2$  {0001} surface to expose, and the small lattice misfit can facilitate the electron transportation through the formation of defects or grain boundaries due to the internal strain, whereas during charge the  $\text{RuO}_2$  monolayer can attract the adsorbed  $\text{Li}_2\text{O}_2$  to move toward its surface to maintain the solid–solid reaction interface. Current work points out that  $\text{RuO}_2$  monolayer is a quite

promising catalytic material for nonaqueous lithium–oxygen batteries. Further experimental and computational explorations on the application of RuO<sub>2</sub> monolayer as well as other cheaper 2D metal oxide materials in nonaqueous lithium–oxygen batteries are expected.

## ■ ASSOCIATED CONTENT

### ● Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcc.6b00014.

Optimized lattice parameters of Li<sub>2</sub>O<sub>2</sub>, RuO<sub>2</sub> monolayer, and rutile RuO<sub>2</sub>. The orientations and terminations of rutile RuO<sub>2</sub> considered for the Wulff construction. Surface energies of different orientations and terminations of rutile RuO<sub>2</sub> under different oxygen chemical potentials. Optimized geometry for oxygen molecular adsorbed onto the surface of RuO<sub>2</sub> monolayer. Optimized geometry for oxygen molecular adsorbed onto rutile RuO<sub>2</sub> {001} surface. Optimized geometry of discharge process happened on Li<sub>2</sub>O<sub>2</sub> {0001} surface with and without RuO<sub>2</sub> monolayer. Energy profile for the discharge process happened on the Li<sub>2</sub>O<sub>2</sub> {0001} surface with and without RuO<sub>2</sub> monolayer. Optimized lattice parameters of Li<sub>2</sub>O<sub>2</sub> before and after adsorption onto RuO<sub>2</sub> monolayer. Illustration for proposed formation mechanism for the assembled thin disc morphology of Li<sub>2</sub>O<sub>2</sub>. The geometries of the interfacial model of rutile RuO<sub>2</sub> {001} surface and Li<sub>2</sub>O<sub>2</sub> {0001} surface before and after optimization. (PDF)

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### Notes

The authors declare no competing financial interest.

## ■ ACKNOWLEDGMENTS

The work described in this paper was fully supported by a grant from the Research Grants Council of the Hong Kong Special Administrative Region, China (Project No. 16213414)

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